



Hybrid Hydro+Micro Models: Status and Outlook

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- 3+1D ideal RFD
- hybrid Hydro+Micro Models
- η/s of a Hadron Gas

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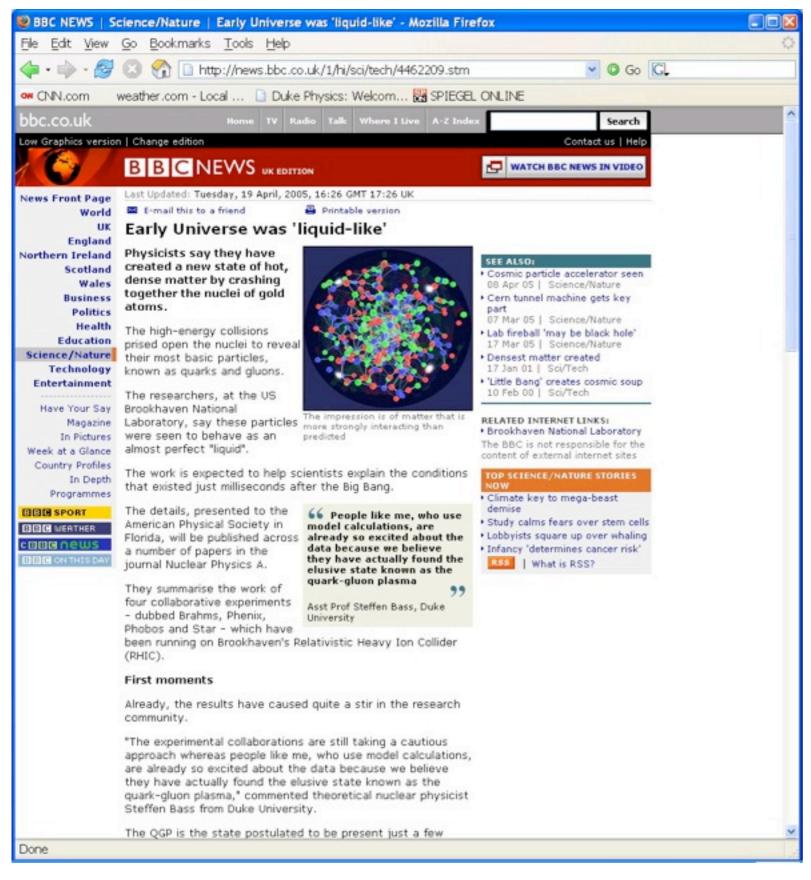


The Case for Hydro & Hydro+Micro



RHIC in the press: Perfect Liquid





- on April 18th, 2005, BNL announced in a press release that RHIC had created a new state of hot and dense matter which behaves like a nearly perfect liquid.
- how does one measure/calculate the properties of an ideal liquid?
- are there any other ideal liquid systems found in nature?



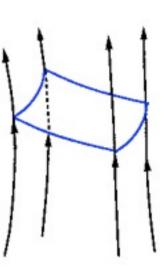




- transport of macroscopic degrees of freedom
- based on conservation laws: $\partial_{\mu}T^{\mu\nu}=0$ $\partial_{\mu}j^{\mu}=0$
- for ideal fluid: $T^{\mu\nu}=(\epsilon+p)$ u^{μ} $u^{\nu}-p$ $g^{\mu\nu}$ and $j_i^{\mu}=\rho_i$ u^{μ}
- Equation of State needed to close system of PDE's: $p=p(T,\rho_i)$
- > connection to Lattice QCD calculation of EoS
- initial conditions (i.e. thermalized QGP) required for calculation
- Hydro assumes local thermal equilibrium, vanishing mean free path

This particular implementation:

- \succ fully 3+1 dimensional, using (τ,x,y,η) coordinates
- Lagrangian Hydrodynamics
 - > coordinates move with entropy-density & baryon-number currents
 - > trace adiabatic path of each volume element





3D-Hydro: Parameters



Initial Conditions:

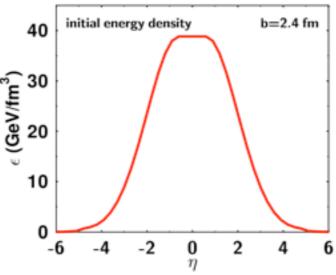
Energy Density:

$$\varepsilon(x, y, \eta) = \varepsilon_{\text{max}} W(x, y; b) H(\eta)$$

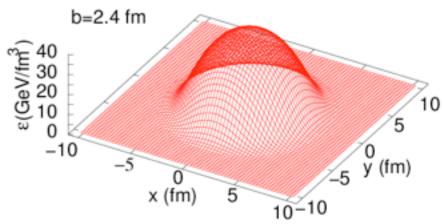
Baryon Number Density:

$$n_B(x, y, \eta) = n_{B\text{max}} W(x, y; b) H(\eta)$$

longitudinal profile:



transverse profile:

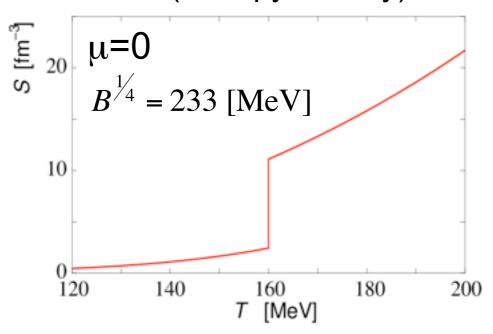


Parameters:

$$\tau_0$$
=0.6 fm/c ϵ_{max} =55 GeV/fm³, n_{Bmax} =0.15 fm⁻³ η_0 =0.5 σ_{η} =1.5

• Initial Flow: $v_L = \eta$ (Bjorken's solution); $v_T = 0$

EOS (entropy density)



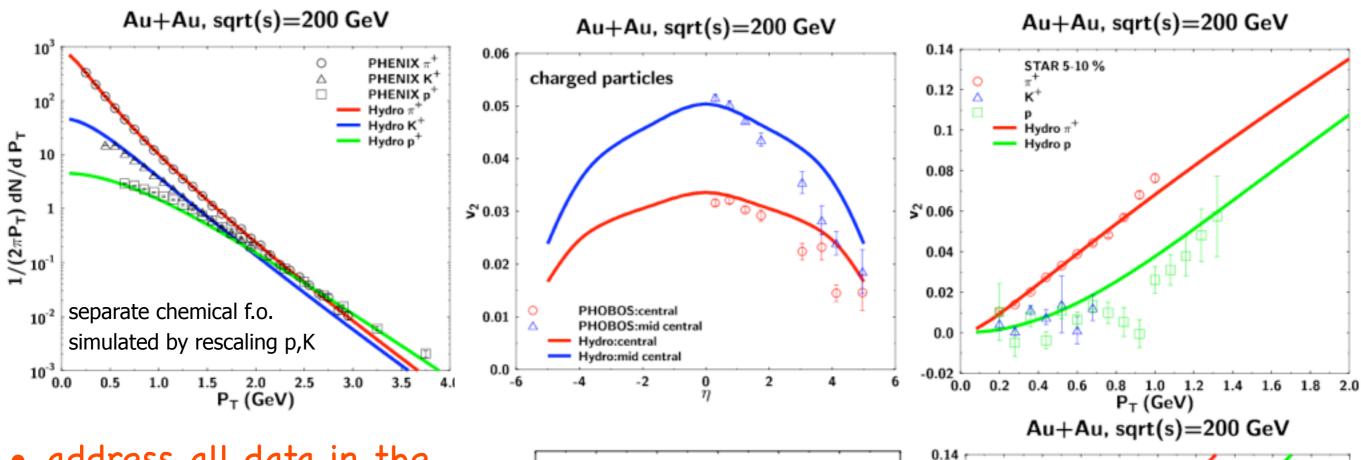
Equation of State:

- Bag Model + excluded volume
- 1st order phase transition (to be replaced by Lattice EoS)

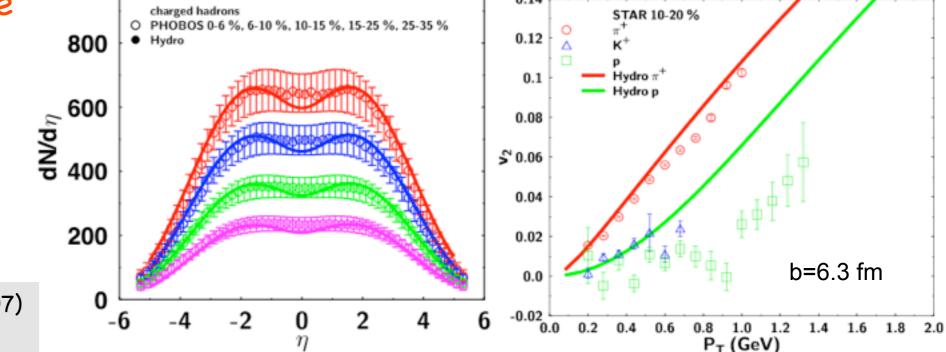


3D-Hydro: Comparison to RHIC





 address all data in the soft sector with one consistent calculation



Nonaka & Bass: PRC75, 014902 (2007)

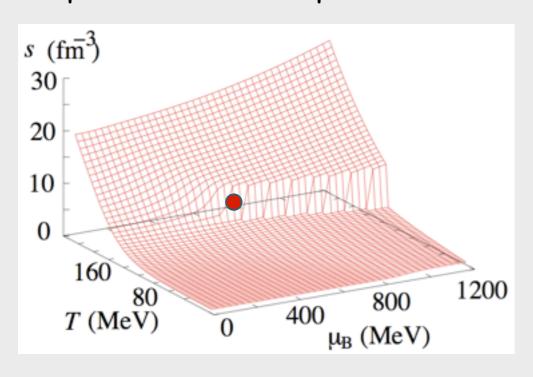
See also Hirano; Kodama et al.

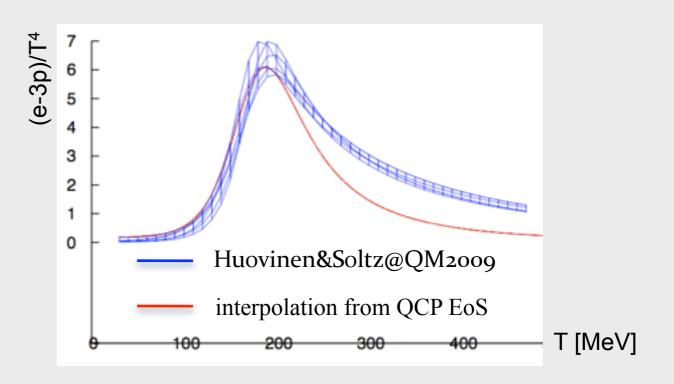


Improved Equation of State

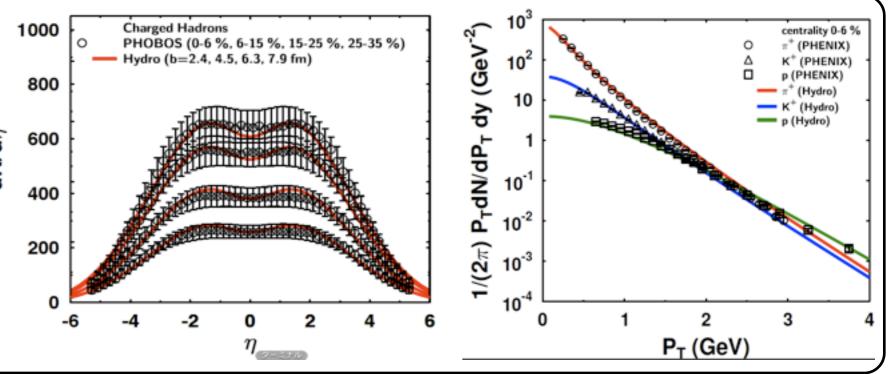


- use EoS with QCD critical point ($T_E=159$ MeV, $\mu_E=550$ MeV)
- interpolate to Lattice parametrization for $\mu_B=0$





- retune initial conditions to $dN_{ch}/d\eta$ and $1/p_{t}$ dN/dp_{t}
- now study effect of EoS on reaction dynamics, jet energy-loss and HBT...
- next item to-do: improved initial conditions based on CGC w/ fluctuations

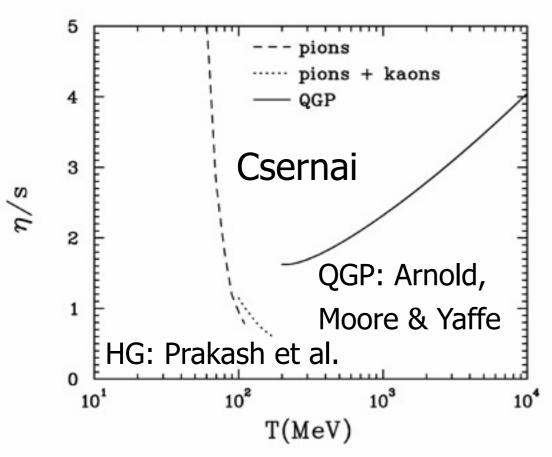




Ideal RFD: Challenges



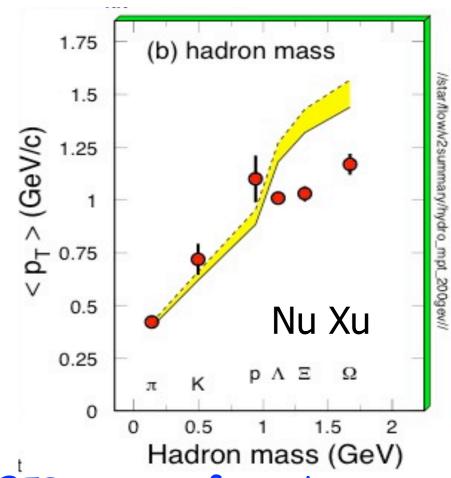
- centrality systematics of v₂ less than perfect
- no flavor dependence of cross-sections
- separation chemical and kinetic freeze-out:
 - normalize spectra by hand
 - PCE: proper normalization, wrong v₂
- off-equilibrium effect!



Viscosity:



- > compatible with AdS/CFT bound of 1/4π
- viscosity will strongly change as function of temperature during collision
- > need to account for viscous corrections, in particular in the hadronic phase





3D-Hydro + Micro Model



full 3-d ideal RFD QGP evolution

Hadronization

Cooper-Frye formula

Monte Carlo

UrQMD

hadronic rescattering

 T_{C}

 T_{SW}

t fm/c



- ideally suited for dense systems
- model early QGP reaction stage
- well defined Equation of State
- parameters:
- initial conditions
- Equation of State

micro. transport (UrQMD)

- no equilibrium assumptions
 - > model break-up stage
 - > calculate freeze-out
 - > includes viscosity in hadronic phase
- parameters:
 - (total/partial) cross sections

matching condition:

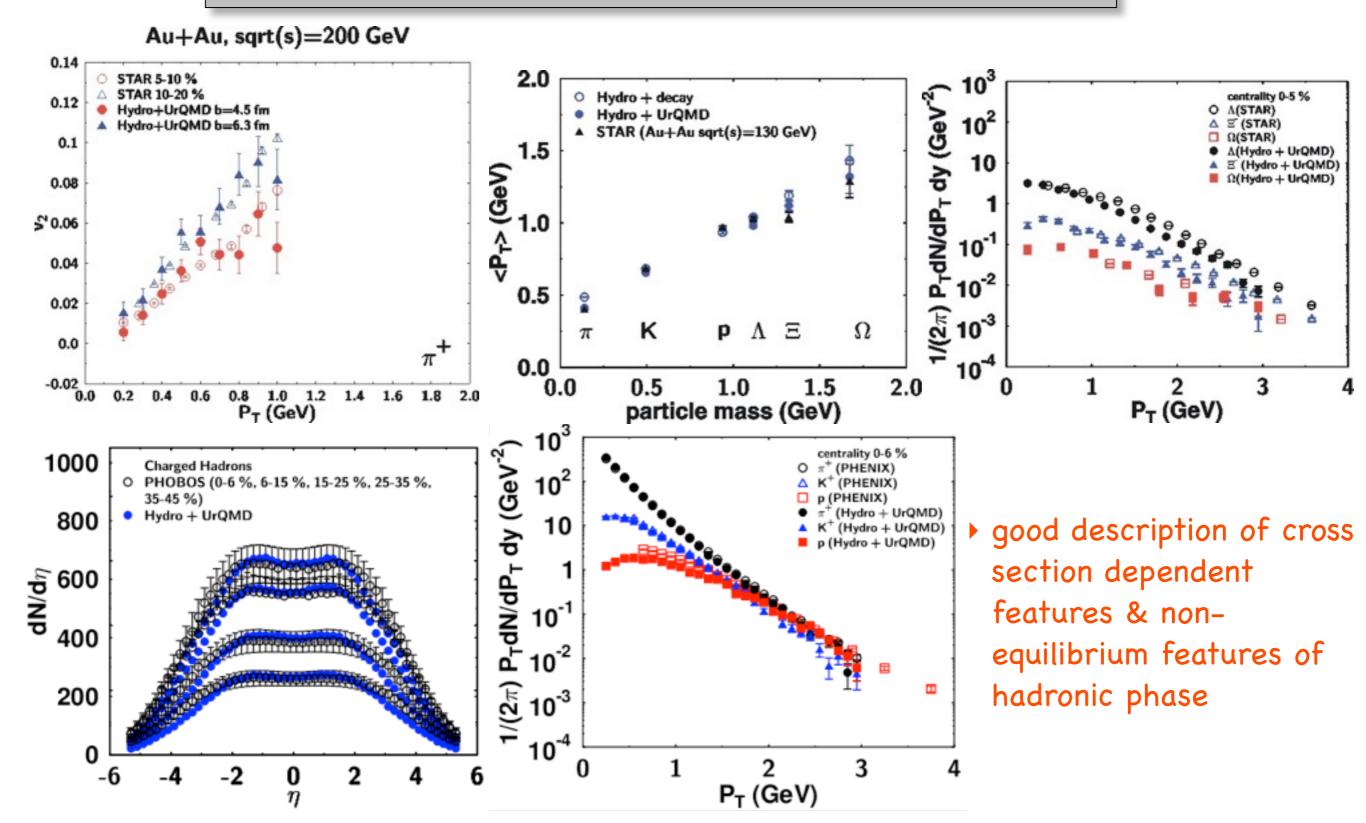
- use same set of hadronic states for EoS as in UrQMD
- \bullet generate hadrons in each cell using local T and μ_{B}
- S.A. Bass & A. Dumitru, Phys. Rev C61 (2000) 064909
- D. Teaney et al, nucl-th/0110037
- T. Hirano et al. Phys. Lett. **B636** (2006) 299
- C. Nonaka & S.A. Bass, Phys. Rev. C75 (2006) 014902















3D-Hydro+Micro: Outlook

- benchmark comparison between different ideal RFD+Micro implementations (TECHQM project / Nonaka to coordinate)
- study sensitivity to EoS
- explore different initial conditions (Glauber vs. CGC beware of extreme scenarios!)
- redo comparison to RHIC data with new EoS & initial conditions
- develop Hydro+Micro converters for vRFD
- develop 3+1D vRFD+Micro as standard model for bulk evolution



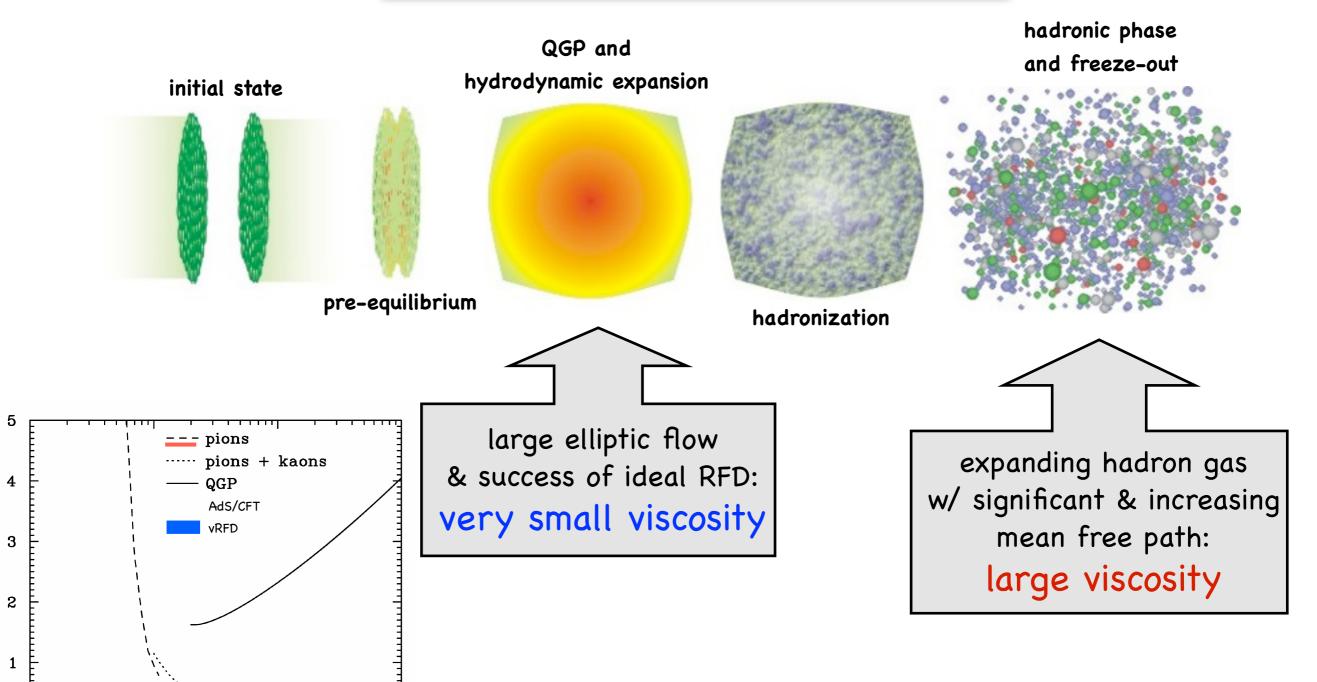


η/s of a Hadron Gas

N. Demir & S.A. Bass: PRL 102, 172302 (2009)



Viscosity at RHIC



L.P. Csernai, J.I. Kapusta & L. McLerran: Phys. Rev. Lett. **97**: 152303 (2006) M. Prakash, M. Prakash, R. Venugopalan & G. Welke: Phys. Rept. **227**, 321 (1993) P. Arnold, G.D. Moore & L.D. Yaffe: JHEP **05**: 051 (2003)

10⁴

10³

10²

T(MeV)

- viscosity of matter @ RHIC changes strongly with time & phase
- how can we learn more about the viscosity of QCD matter?



Microscopic Transport: η/s of a Hadron Gas



• for particles in a fixed volume, the stress energy tensor discretizes

$$\pi^{xy} = \frac{1}{V} \sum_{j=1}^{N_{\text{part}}} \frac{p^{x}(j)p^{y}(j)}{p^{0}(j)}$$

• and the Green-Kubo formula reads:

$$\eta = \frac{V}{T} \int_0^\infty dt \, \langle \pi^{xy}(0) \, \pi^{xy}(t) \rangle$$

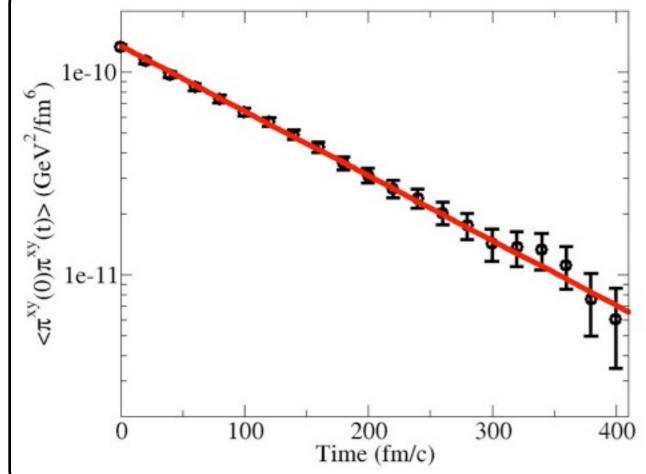
Entropy:

• extract thermodynamic quantities via:

$$P = \frac{1}{3V} \sum_{j=1}^{N_{\text{part}}} \frac{|\vec{p}|^2(j)}{p^0(j)} \quad \epsilon = \frac{1}{V} \sum_{j=1}^{N_{\text{part}}} p^0(j)$$

use Gibbs relation (with chem. pot. extratced via SM)

$$s_{\text{Gibbs}} = \left(\frac{\epsilon + P - \mu_i \rho_i}{T}\right)$$



- evaluating the correlator numerically, e.g. in UrQMD one empirically finds an exponential decay as function of time
- using the following ansatz, one can extract the relaxation time T_{π} :

$$\langle \pi^{xy}(0) \, \pi^{xy}(t) \rangle \propto \exp\left(-\frac{t}{\tau_{\pi}}\right)$$

 the shear viscosity then can be calculated from known/extracted quantities:

$$\eta = \frac{\tau_{\pi}}{T} \left\langle \pi^{xy}(0)^2 \right\rangle$$

A. Muronga: Phys. Rev. C69: 044901, 2004

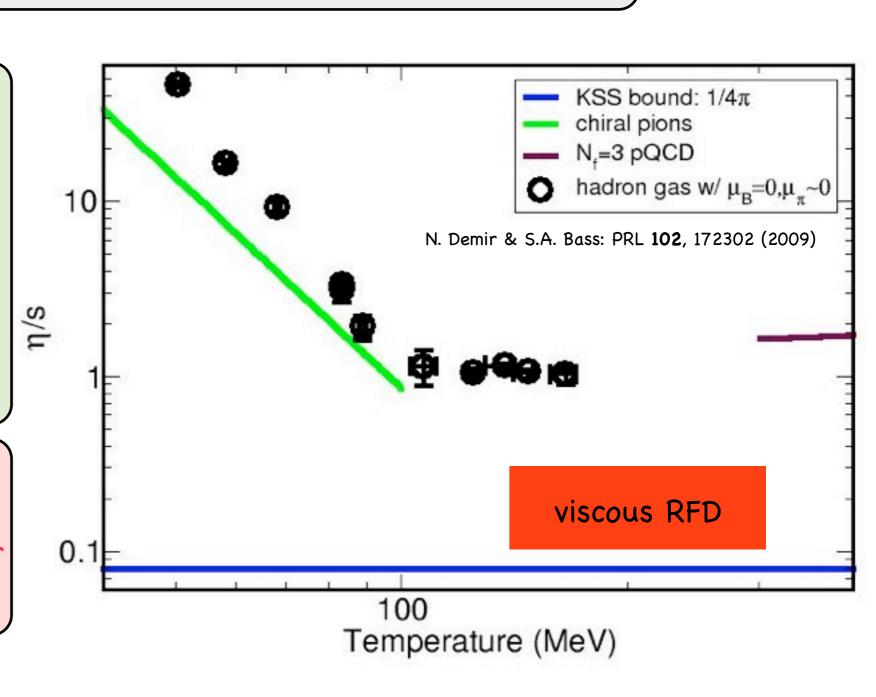


η/s of a Hadron Gas



first reliable calculation of of η/s for a full hadron gas including baryons and anti-baryons

- low temperature trend qualitatively confirms chiral pion calculation
- above T=100 MeV: η/s≈1
 remains roughly constant
- η/s is a factor of 3-5 above range required by viscous RFD analysis!
- breakdown of vRFD in the hadronic phase?
- what are the consequences for η/s in the deconfined phase?





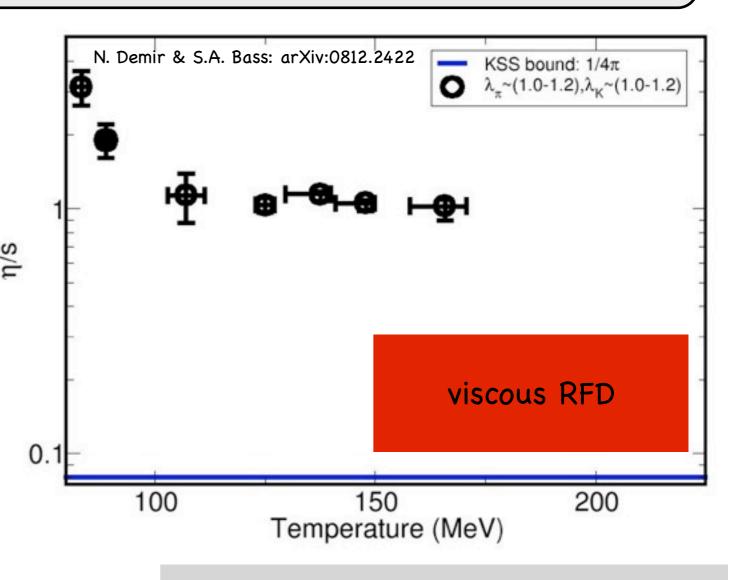
Dynamic Systems: η/s at non-unit fugacities



- freeze-out temperature required by RFD to reproduce spectral shapes: ~110 MeV
- temperature extracted from Statistical Model fits to hadron yields/ratios: ~160 MeV
- ▶ separation of chemical and kinetic freeze-out in the hadronic phase!
- picture confirmed by hybrid hydro+micro calculations
- ▶ off-equilibrium effect implies non-unit species-dependent fugacities in RFD

box calculations w/ non-unit fugacities:

- initialize matter w/ equilibrium distributions, but off-equilibrium yields, corresponding to desired fugacities
- perform viscosity measurement before system relaxes into equilibrium
- verify fugacities at time of measurement w/ statistical model analysis
- non-unit fugacities reduce η/s by a factor of two to η/s≈0.5
- η/s still above value required for viscous RFD fit to data
- η/s needs to be significantly lower in deconfined phase for vRFD to reproduce elliptic flow data!



T. Hirano & K. Tsuda: Nucl. Phys. A715, 821 (2003)

P.F. Kolb & R. Rapp: Phys. Rev. C67, 044903 (2003)



Conclusions and Outlook

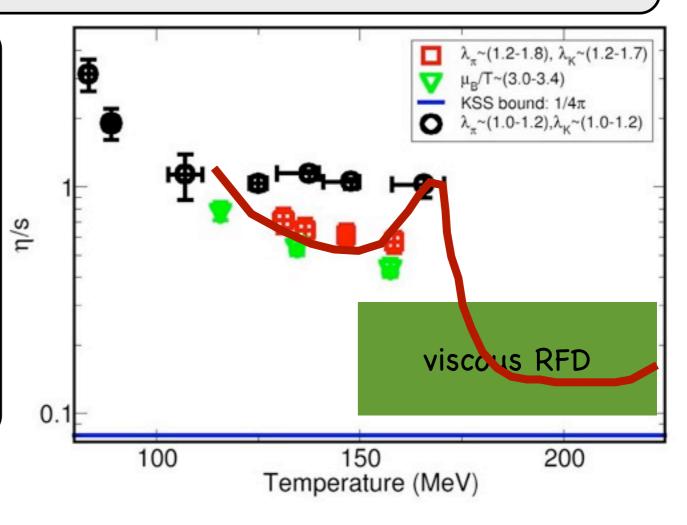


Hybrid Hydro+Micro models:

- 3+1D Hydro + Micro models have been very successful in describing the bulk properties of hot & dense QCD matter created at RHIC
- the microscopic treatment of the hadronic phase does not only address viscous effects, but also the inherent off-equilibrium evolution of the system during its break-up stage
- in the future, 3+1D vRFD + Micro models should be pursued, combining the best possible description of the low viscosity deconfined phase with the optimum description of the hadronic phase

Viscosity of QCD matter:

- need to parametrize η/s as function of T for vRFD calculations
- trajectory of η/s in a heavy-ion collision as a function of temperature may have complicated shape
- calculation of hadronic η/s will help to constrain η/s in the deconfined phase





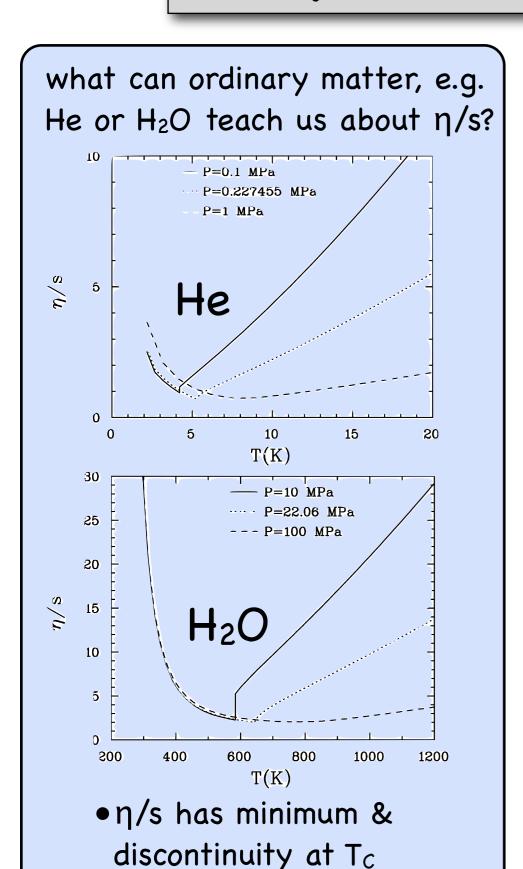


The End



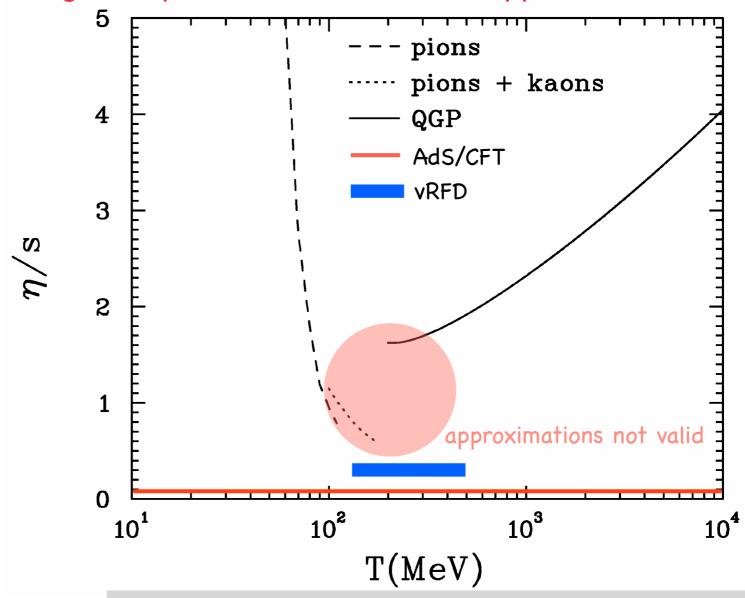






temperature dependence of η/s in QCD can be estimated in low- and high-temperature limit:

- low temperature: chiral pions
- high temperature: QGP in HTL approximation



L.P. Csernai, J.I. Kapusta & L. McLerran: Phys. Rev. Lett. **97**: 152303 (2006)
M. Prakash, M. Prakash, R. Venugopalan & G. Welke: Phys. Rept. **227**, 321 (1993)
P. Arnold, G.D. Moore & L.D. Yaffe: JHEP **05**: 051 (2003)



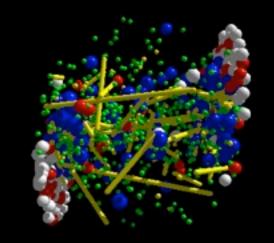
Hadronic Matter: UrQMD



- •elementary degrees of freedom: hadrons, const. (di)quarks
- •classical trajectories in phase-space (relativistic kinematics): evolution of phase-space distribution via Boltzmann Equation:

$$\left[\frac{\partial}{\partial t} + \frac{\vec{p}}{m} \nabla_r\right] f^1 = \mathcal{C}_{\rm coll}$$
 with $\mathcal{C}_{\rm coll} = N \int \sigma \mathrm{d}\Omega \int \mathrm{d}\vec{p}_2 \, |\vec{v}_1 - \vec{v}_2| \, [f_1(\vec{p}_1') f_1(\vec{p}_2') - f_1(\vec{p}_1) f_1(\vec{p}_2)]$

- initial high energy phase of the reaction is modeled via the excitation and fragmentation of strings
- \bullet 55 baryon- and 32 meson species, among those 25 N*, Δ * resonances and 29 hyperon/hyperon resonance species
- full baryon-antibaryon and isospin symmetry



main physics input and parameters:

- cross sections: total and partial cross sections, angular distributions
- resonance parameters: total and partial decay widths
- string fragmentation scheme: fragmentation functions, formation time

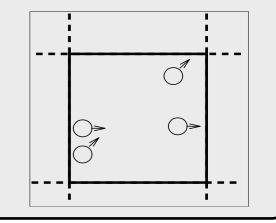


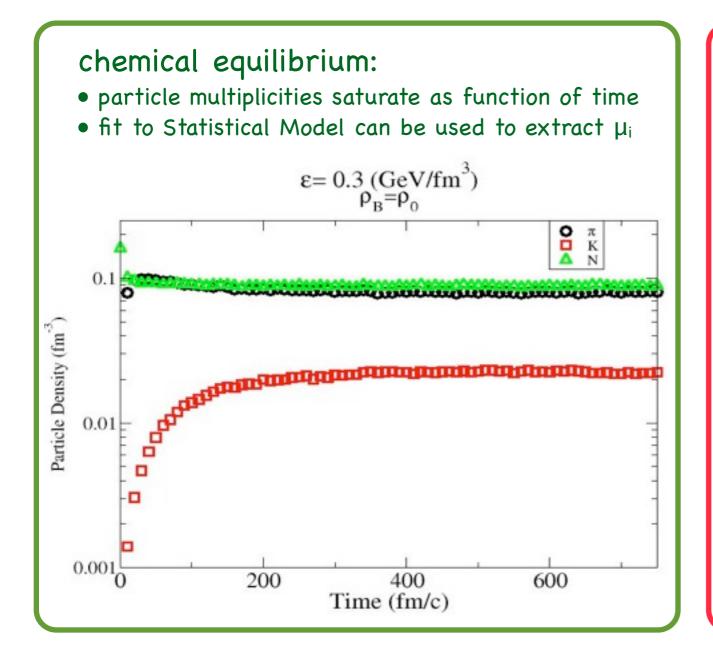
Infinite Matter Calculations

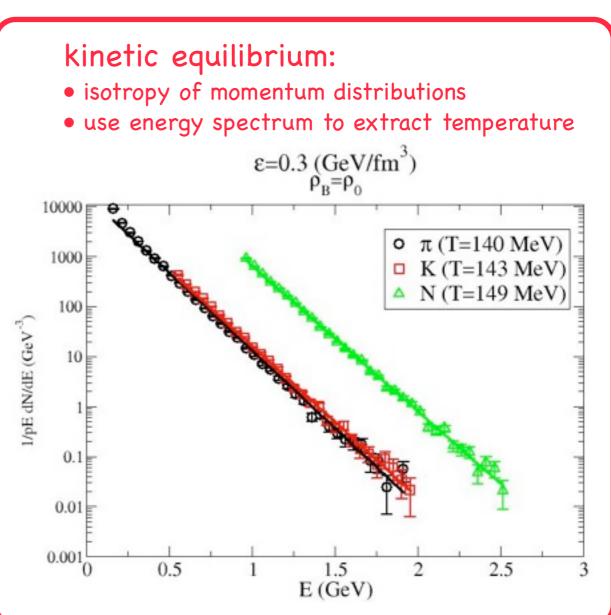


Strategy: confine UrQMD to box with periodic boundary conditions

- system will evolve into equilibrium state (no freeze-out occurs)
- need to disable multi-body processes to maintain detailed balance









Shear Viscosity: Linear Transport Equation & Green - Kubo Formalism



Mechanical definition of shear viscosity:

•application of a shear force to a system gives rise to a non-zero value of the xy-component of the pressure tensor P_{xy} . P_{xy} is then related to the velocity flow field via the shear viscosity coefficient η :

 $P_{xy} = -\eta \frac{\partial v_x}{\partial y}$

- a similar linear transport equation can be defined for other transport coefficients: thermal conductivity, diffusion ...
- •using linear-response theory, the Green-Kubo relations for the shear viscosity can be derived, expressing η as an integral of an near-equilibrium time correlation function of the stress-energy tensor:

$$\eta = \frac{1}{T} \int d^3r \int_0^\infty dt \left\langle \pi^{xy}(\vec{0}, 0) \pi^{xy}(\vec{r}, t) \right\rangle_{\text{equil}}$$

with the stress-energy tensor: $\pi^{\mu\nu}(\vec{r},t)=\int d^3p \frac{p^\mu p^\nu}{p^0} f(x,p)$



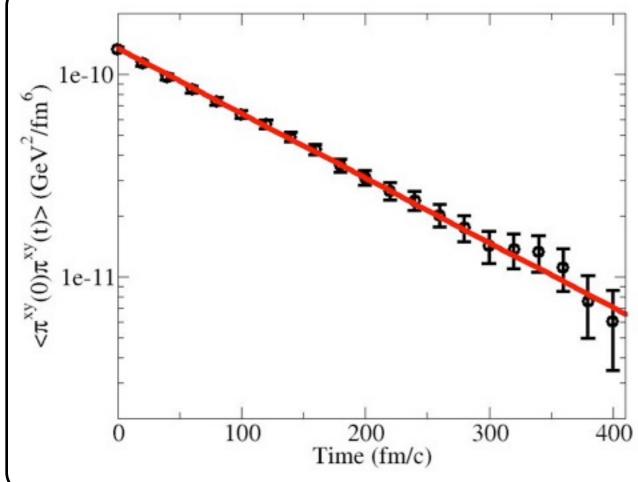
Kubo Formalism in Microscopic Transport



• for a set of discrete particles in a fixed volume, the stress energy tensor discretizes

$$\text{from} \ \ \pi^{\mu\nu}(\vec{r},t) = \int d^3p \frac{p^\mu p^\nu}{p^0} f(x,p) \quad \text{ to } \pi^{xy} = \frac{1}{V} \sum_{i=1}^{N_{\mathrm{part}}} \frac{p^x(j) p^y(j)}{p^0(j)}$$

ullet and the Green-Kubo formula reads: $m{\eta}=rac{V}{T}\int_0^\infty dt \ \langle \pi^{xy}(0)\,\pi^{xy}(t)
angle$



- evaluating the correlator numerically, e.g. in UrQMD one empirically finds an exponential decay as function of time
- •using the following ansatz, one can extract the relaxation time T_{π} :

$$\langle \pi^{xy}(0) \, \pi^{xy}(t) \rangle \propto \exp\left(-\frac{t}{\tau_{\pi}}\right)$$

 the shear viscosity then can be calculated from known/extracted quantities:

$$\eta = \frac{1}{\tau_{\pi}} \frac{V}{T} \left\langle \pi^{xy}(0)^2 \right\rangle$$

A. Muronga: Phys. Rev. C69: 044901, 2004



Entropy in Microscopic Transport Models



The extraction of entropy from microscopic transport models is non-trivial:

- •use two independent methods to ensure accuracy
- \bullet thermodynamic quantities which can be extracted directly from box-calculation are: pressure p, energy-density ϵ , particle number N_i , temperature T and volume V

$$P = \frac{1}{3V} \sum_{j=1}^{N_{\text{part}}} \frac{|\vec{p}|^2(j)}{p^0(j)} \qquad \epsilon = \frac{1}{V} \sum_{j=1}^{N_{\text{part}}} p^0(j)$$

Method #1: Gibbs entropy

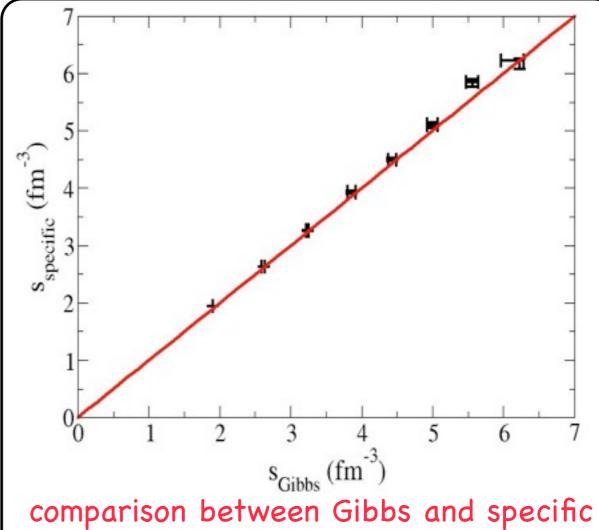
- extract chemical potential(s) from SM
- use Gibbs-relation for entropy:

$$s_{\mathrm{Gibbs}} = \left(\frac{\epsilon + P - \mu_i \rho_i}{T}\right)$$

Method #2: specific entropy

• sum over specific entropies of all hadron species, which can be calculated as functions of m/T and μ_B/T :

$$s_{ ext{specific}} = rac{1}{V} \sum_{i}^{N_{spec}} \left(rac{s}{n}
ight)_{i} N_{i}$$



comparison between Gibbs and specific entropy shows excellent agreement!







The consistency of the entropy extraction can be verified via a scaling law with the speed of sound of the system: 1

 $s \sim T^{\frac{1}{c_s^2}}$

Step #1: determine speed of sound c_s , using pressure and energy-density:

$$c_s^2 = \left(\frac{\partial P}{\partial \epsilon}\right)$$

• analysis yields $c_s^2=0.18$

Step #2: plot Gibbs entropy vs. temperature, using the scaling law

scaling law is well reproduced

